# A Modified Formula for Microwave Measurement of Dielectric Loss Using a Closed Cylindrical Cavity Dielectric Resonator

## Liangzu Cao<sup>1, 2, \*</sup> and Daming Cao<sup>3</sup>

Abstract—This paper provides a modified formula for calculating dielectric loss of dielectric resonator of working in  $TE_{01\delta}$  mode in closed cavity. The measurement system is divided into six regions. Based on analyzing the formula of loss tangent published in literatures, a quality factor of a substrate is created, and a modified formula is proposed. Validating the modified formula, with three substrates as supports, the frequencies and unloaded quality factors of dielectric resonators made of two sorts of dielectric materials with permittivity 38 and 75 respectively are measured using a closed cavity method. The measured results are compared with those obtained by other well-known formulas and show a good agreement with the result given by the parallel plate method.

### 1. INTRODUCTION

Dielectric materials continue to play a very important role in the microwave communication systems. These materials are key in realization of low-loss temperature-stable resonators, filters for satellite and broadcasting equipment, as well as in many other microwave devices [1-3]. The microwave measurements of dielectric properties have recently attracted more and more attention. There are many measurement techniques [4,5], such as Hakki-Coleman end shorted method [6,7],  $TE_{01\delta}$  mode dielectric resonator method [8,9], Whispering gallery mode resonators [10–12] and Cavity perturbation method [13]. Hakki-Coleman end shorted method, called parallel plate method, is the most famous resonance technique. However, the method has very strict requirement, for example, measured dielectric sample being standard cylinder with restrict ratio of diameter to height, surface resistance of metal plates to be calibrated before measurement [14, 15]. TE<sub>01 $\delta$ </sub> mode dielectric resonator method, called closed cavity resonance method, is a simple and novel measurement technique. Putting the dielectric resonator inside a closed metal cavity can completely eliminate the effect of field radiation. Several different theoretical and experimental analyses have been reported in the literatures [16–19]. Accurate prediction of resonant frequency was given in their papers; however the field expressions are usually quite complicated and imperfect. A modified field model has been proposed by Sheen [20, 21], and the dielectric constant was measured with reasonable accuracy; however the effect of a support on loss tangent was neglected. Experiments showed that the support affected the unloaded quality factor [22], but the relationship between the properties of a support and loss tangent of the dielectric resonator has not been studied vet.

This paper adopts the modified field model. A modified formula for calculating the loss tangent of dielectric resonator has been proposed in this paper. To validate the formula, ceramic resonators are measured using the closed cavity method. The results are then compared with those obtained by the open cavity method.

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<sup>\*</sup> Corresponding author: Liangzu Cao (clz4233@aliyun.com).

<sup>&</sup>lt;sup>1</sup> School of Electric and Optical Engineering, Nanjing University of Science & Technology, Nanjing 210014, China. <sup>2</sup> School of Mechanical and Electric Engineering, Jingdezhen Ceramics Institute, Jingdezhen 333403, China. <sup>3</sup> School of Information Science and Engineering, Southeast University, Nanjing 211100, China.

#### 2. DIELECTRIC LOSS MEASUREMENT

Figure 1 shows the structure model of a closed cylindrical cavity dielectric resonator. The cavity is divided into six regions, where region 1 represents cylindrical dielectric resonator; regions 2 and 4 are substrate; regions 3, 5 and 6 are filled with air.



Figure 1. Vertical section view of a closed cavity resonator [20].

When the resonance mode of the dielectric resonator is  $TE_{01\delta}$ , the magnetic field expressions of the z component in each region,  $H_{zi}$   $(i = 1 \sim 6)$ , are given by Sheen [20]. Other components of  $TE_{01\delta}$  in each region can be derived from literature [23].

There are various formulas for calculating dielectric loss of a resonator, according to the literatures. For example, a few are shown below:

Ni [24]: 
$$\tan \delta = \frac{1}{Q_d} = \frac{1}{Q_u} - \frac{1}{Q_c}$$
 (1)

Gu [25]: 
$$\tan \delta = \frac{1}{Q_d} = \frac{1}{Q_u} - \frac{1}{Q_c} - \frac{1}{Q_r}$$
 (2)

Zhou [23]: 
$$\tan \delta = (1+A)\frac{1}{Q_u} - B$$
 (3)

Sheen [21]: 
$$\tan \delta = \frac{W_a + W_d}{W_d} \left( \frac{1}{Q_u} - \frac{1}{Q_c} \right)$$
 (4)

where  $Q_u$  is the unloaded quality factor of the whole system, and  $Q_d$ ,  $Q_c$  and  $Q_r$  represent the quality factor resulting from the dielectric loss, conducing loss and radiation loss, respectively.  $W_d$  and  $W_a$ represent the energy stored in the dielectric and surrounding regions. A and B are the ratios related with the stored energy and the dissipated power, respectively [23].

These formulas are valid under certain conditions, which are beyond the scope of this paper.

A modified formula is proposed in this paper, written as:

$$\tan \delta = \frac{W_a + W_d + W_{sub}}{W_d} \cdot \frac{1}{Q_u} - \frac{P_{wall}}{\omega_0 W_d} - \frac{\tan \delta_{sub} \cdot W_{sub}}{W_d}$$
(5)

where  $W_{sub}$  and  $W_a$  represent the energy stored in the substrate used as a support and the air, respectively, and  $P_{wall}$  denotes the dissipated power in the conducting wall and  $\tan \delta_{sub}$  the dielectric loss of the substrate. Equation (5) can be equivalently written as (6) by mathematical manipulations:

$$\tan \delta = A \left( \frac{1}{Q_u} - \frac{1}{Q_c} - \frac{1}{Q_{sub}} \right) \tag{6}$$

where  $A = \frac{W_a + W_d + W_{sub}}{W_d}$ ,  $Q_c = \omega_0 \frac{W_a + W_d + W_{sub}}{P_{wall}}$ ,  $Q_{sub} = \frac{W_a + W_d + W_{sub}}{\tan \delta_{sub} \cdot W_{sub}}$ . The values of  $W_{sub}$ ,  $W_a$ , and  $W_d$  are calculated by,

$$W_d = \frac{1}{2} \varepsilon_r \varepsilon_0 \iiint |E_{\varphi 1}|^2 \, dV \tag{7a}$$

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$$W_a = \frac{1}{2} \varepsilon_0 \sum_{i=3,5,6} \iiint |E_{\varphi i}|^2 \, dV \tag{7b}$$

$$W_{sub} = \frac{1}{2} \varepsilon_{r1} \varepsilon_0 \sum_{i=2,4} \iiint |E_{\varphi i}|^2 dV$$
(7c)

The value of  $P_{wall}$  is calculated by [9],

$$P_{wall} = \frac{1}{2} R_s \sum_{i=2}^{5} \iint |H_{ri}|^2 \, dS + \frac{1}{2} R_s \sum_{i=4}^{6} \iint |H_{zi}|^2 \, dS \tag{8}$$

where  $\varepsilon_r$  and  $\varepsilon_{r1}$  are relative dielectric constants of resonator and the substrate, respectively;  $E_{\phi i}$  and  $H_{ri}$  denote circumferential electric field components and radial magnetic field components in each region;  $R_s$  is the surface resistance of the metal wall.

Derivation of (5) is described in the Appendix.

## 3. EXPERIMENTS AND DISCUSS

Two resonators are made of ceramics, whose composites are  $(Zr_{0.8}Sn_{0.2})TiO_4$  and BaO-TiO<sub>2</sub>-Sm<sub>2</sub>O<sub>3</sub>, and whose dielectric constants are about 38 and 75, respectively. The cavity is made from aluminium which has the conductivity of  $3.54 \times 10^7$  S/m and an inside diameter/height ratio = 35 mm/17.5 mm. There are three substrates made of plastic, alumina ceramics and Teflon glass textile board (TFGT), respectively. Their dielectric constants are 3.68, 6.75 and 2.65; their dielectric losses are 0.043, 0.0018 and 0.012; their heights are 5.44 mm, 4.32 mm and 1.26 mm; respectively. The resonant frequency and the unload quality factor of resonators are measured using Vector Networks Analyzer (Model: Agilent E5071B). Figure 2 illustrates the experimental setup used for the measurement.



Figure 2. Measurement setup of the closed cylindrical cavity dielectric resonator.

Table 1 shows the measurement results in comparison with the results by the formulas in the literatures. The loss tangent measurements of the same sample from (5) show very close results among the substrates made of plastic, ceramics and TFGT. The results are very close to those measured from the parallel plate method, compared to the results by other formulas. This confirms the accuracy of (5) proposed in this paper.

Loss tangents are same for (3) and (4), because (3) can be changed into (4) by substituting A and B with their expressions and simplifying it. Equations (3) or (4) does not suit for calculating dielectric loss with high dielectric loss substrate as support, because the quality factor of the substrate is much less than that of the cavity. For example,  $Q_{sub}/Q_c$  of Sample#1 using plastic as support is  $6.51 \times 10^{-3}$ , shown in Table 2. The higher is the loss tangent of the substrate, the larger is the relative error of loss tangent calculated using (3). As for TFGT, the relative errors are more than 50%, even 90% for Sample#1.

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	Substrate	Freq.	$Q_u$	$ an \delta$	$ an \delta$	$ an \delta$	$ an \delta$	$\tan \delta$ from
		(GHz)		from $(5)$	from $(3)$	from $(4)$	from $(1)$	Parallel
								$plate^*$
Sample#1	Plastic	4.063	1120	$8.33 \times 10^{-5}$	$9.84 \times 10^{-4}$	$9.84 \times 10^{-4}$	$8.84 \times 10^{-4}$	
$D = 12.84 \mathrm{mm}$	Ceramics	4.023	6800	$8.72 \times 10^{-5}$	$1.54\! imes\!10^{-4}$	$1.54\! imes\!10^{-4}$	$1.37\!\times\!10^{-4}$	$7.856\!\times\!10^{-5}$
$L\!=\!5.50\mathrm{mm}$	TFGT	4.308	6080	$8.70 \times 10^{-5}$	$1.74 \times 10^{-4}$	$1.74 \times 10^{-4}$	$1.59\!\times\!10^{-4}$	
Sample#2	Plastic	2.868	1340	$3.08 \times 10^{-4}$	$7.74 \times 10^{-4}$	$7.74 \times 10^{-4}$	$7.40 \times 10^{-4}$	
$D = 13.02 \mathrm{mm}$	ceramics	2.862	2860	$3.26\! imes\!10^{-4}$	$3.62\! imes\!10^{-4}$	$3.62\! imes\!10^{-4}$	$3.43 \times 10^{-4}$	$3.5\!\times\!10^{-4}$
$L\!=\!5.50\mathrm{mm}$	TFGT	3.043	2400	$3.29\! imes\!10^{-4}$	$3.43\!\times\!10^{-4}$	$3.43\!\times\!10^{-4}$	$3.31\!\times\!10^{-4}$	

 Table 1. Dielectric properties as measured using various formulas.

\*Loss tangent at certain frequency is derived from the product of frequency and the reciprocal of loss tangent.

 Table 2. Calculated parameters.

	Substrate	$Q_c$	$Q_{sub}$	$Q_{sub}/Q_c$	$(W_a + W_{sub})/W_t^{**}$	$W_d/W_t$	A
Sample#1	Plastic	$1.90 \times 10^5$	1237	$6.51 \times 10^{-3}$	0.1019	0.8981	1.1088
$D\!=\!12.84\mathrm{mm}$	Ceramics	$1.00\!\times\!10^5$	$1.69\!\times\!10^4$	0.169	0.1078	0.8922	1.1210
$L\!=\!5.50\mathrm{mm}$	TFGT	$1.60\!\times\!10^4$	$4.40\!\times\!10^4$	2.75	0.0847	0.9153	1.0920
Sample #2	Plastic	$1.67\!\times\!10^5$	2326	$1.39\! imes\!10^{-2}$	0.0441	0.9559	1.0460
$D = 13.02 \mathrm{mm}$	ceramics	$1.59\!\times\!10^5$	$2.92\!\times\!10^4$	0.184	0.0518	0.9482	1.0546
$L\!=\!5.50\mathrm{mm}$	TFGT	$1.17\!\times\!10^5$	$7.70\!\times\!10^4$	0.658	0.0344	0.9656	1.0356

 $^{**}W_t = W_a + W_{sub} + W_d$ 

It is strange that  $Q_{sub}/Q_c$  of Sample#1 using TFGT as support is 2.75, and  $Q_c$  value is  $1.60 \times 10^4$ , listed in Table 2, much smaller than others. The reason is that the thickness of TFGT is 1.26 mm, and dielectric constant of Sample#1 is lower, about 38, which make  $Q_c$  value smaller. It is inferred that the supports with high dielectric loss have much more effect on the measured Q value of dielectric resonators with low dielectric constant, and the effect of support cannot be ignored and should be taken into account.

#### 4. CONCLUSIONS

The unloaded quality factor of a closed cavity resonator is dependent on the substrate used as a support. The larger is the loss tangent of the substrate, the smaller is the unloaded quality factor of the system. The modified formula proposed in this paper shows that the dielectric loss of a resonator can be calculated by subtracting the conducting loss  $(1/Q_c)$  and the substrate dissipated loss  $(1/Q_{sub})$  from the reciprocal of the unloaded quality factor  $(1/Q_u)$ . The accuracy of the proposed formula is confirmed by the results from the experiments. Compared to other formulas in the literature, the results by the proposed formula show a much closer match with those by the parallel plate method, and the results are robust even with different materials for the substrate.

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#### APPENDIX A.

Unloaded quality factor  $(Q_u)$  of dielectric resonator located in a closed cavity is defined as [2]

$$Q_u = 2\pi \frac{\text{maximum energy stored per cycle}}{\text{Average energy dissipated per cycle}} = \frac{\omega_0 W_t}{p_t}$$
(A1)

where  $\omega_0$  is angle frequency,  $W_t$  the total energy stored in all media, and  $P_t$  the total power dissipated by all substances.

$$W_t = W_a + W_d + W_{sub} \tag{A2}$$

$$p_t = p_a + p_d + p_{sub} + p_{wall} \tag{A3}$$

where  $P_d$ ,  $P_{sub}$ ,  $P_{wall}$  and  $P_a$  represent the power dissipated in the dielectric, substrate, conductor and the air respectively, and  $W_a$ ,  $W_d$  and  $W_{sub}$  are mentioned in Section 2.

Loss tangents of dielectric resonator and substrate are written as

$$\tan \delta_{sub} = \frac{p_{sub}}{\omega_0 W_{sub}} \tag{A4}$$

$$\tan \delta_d = \frac{p_d}{\omega_0 W_d} \tag{A5}$$

By adding (A1)–(A4) into (A5), (A5) is rewritten as

$$\tan \delta_d = \frac{W_t}{W_d Q_u} - \frac{p_a}{\omega_0 W_d} - \frac{\tan \delta_{sub} W_{sub}}{W_d} - \frac{p_{wall}}{\omega_0 W_d}$$
(A6)

When we ignore the power dissipated in the air, i.e.,  $p_a = 0$ , (A7) can be obtained.

$$\tan \delta_d = \frac{W_t}{W_d Q_u} - \frac{\tan \delta_{sub} W_{sub}}{W_d} - \frac{p_{wall}}{\omega_0 W_d} \tag{A7}$$

#### REFERENCES

- Fiedziuszko, S. J., I. C. Hunter, T. Itoh, et al., "Dielectric materials, devices, and circuits," *IEEE Trans. Microwave Theory Tech.*, Vol. 50, No. 3, 706–720, Mar. 2002.
- 2. Sebastian, M. T., Dielectric Materials for Wireless Communications, Elsevier Ltd, USA, 2008.
- Kajfezz, D. and P. Guillon, *Dielectric Resonators*, Noble Publishing Corporation, Tucker, Georgia, USA, 1998.
- 4. Sheen, J., "Comparisons of microwave dielectric property measurements by transmission/reflection techniques and resonance techniques," *Meas. Sci. Technol.*, Vol. 20, 1–12, 2009.
- 5. Krupka, J., "Precise measurements of the complex permittivity of dielectric materials at microwave frequencies," *Materials Chemistry and Physics*, Vol. 79, 195–198, 2003.
- Hakki, B. W. and P. D. Coleman, "A dielectric resonator method of measuring inductive capacities in the millimeter range," *IRE Trans. Microwave Theory Tech.*, Vol. 8, No. 7, 402–410, 1960.
- Courtney, W. E., "Analysis and evaluation of a method of measuring complex permittivity and permeability of microwave materials," *IEEE Trans. Microwave Theory Tech.*, Vol. 18, 476–485, 1970.
- Krupka, J., "Frequency domain complex permittivity measurements at microwave frequencies," Meas. Sci. Technol., Vol. 16, R1–R16, 2005.
- Krupka, J., K. Derzakowski, B. Riddle, and J. B. Jarvis, "A dielectric resonator for measurements of complex permittivity of low loss materials as a function of temperature," *Meas. Sci. Technol.*, Vol. 9, 1751–1756, 1998.
- Cros, D. and P. Guillon, "Whispering gallery dielectric resonator modes for W-band devices," *IEEE Microw. Theory Tech.*, Vol. 38, 1667–1674, 1990.
- 11. Krupka, J., D. Cros, M. Aubourg, and P. Guillon, "Study of whispering gallery modes in anisotropic single crystal dielectric resonators," *IEEE Trans. Microwave Theory Tech.*, Vol. 42, 56–61, 1994.

- 12. Krupka, J., K. Derzakowski, A. Abramowicz, M. E. Tobar, and R. G. Geyer, "Whispering gallery modes for complex permittivity measurements of ultra low loss dielectric materials," *IEEE Microw. Theory Tech.*, Vol. 47, 752–759, 1999.
- Prakash, A., J. K. Vaid, and A. Mansingh, "Measurement of dielectric parameters at microwave frequencies by cavity perturbation technique," *IEEE Trans. Microwave Theory Tech.*, Vol. 27, 791–795, 1979.
- Kobsyashi, Y. and M. Katon, "Microwave measurement of dielectric properties of low-loss materials by the dielectric rod resonator method," *IEEE Trans. Microwave Theory Tech.*, Vol. 33, No. 7, 586– 592, 1985.
- 15. Cao, L. and D. Cao, "Fast measurement of complex permittivity of microwave dielectric materials using parallel short-circuit plate method," *Journal of Ceramics*, Vol. 33, No. 3, 80–84, 2012, in Chinese.
- 16. Bonetti, R. and A. Atia, "Design of cylindrical dielectric resonators in inhomogeneous media," *IEEE Trans. Microwave Theory Tech.*, Vol. 29, No. 4, 323–326, 1981.
- 17. Itoh, T. and R. Rudokas, "New method for computing the resonant frequency of dielectric resonator," *IEEE Trans. Microwave Theory Tech.*, Vol. 25, 52–54, 1977.
- Maystre, D., P. Vincent, and J. C. Mage, "Theoretical and experimental study of the resonant frequency of a cylindrical dielectric resonator," *IEEE Trans. Microwave Theory Tech.*, Vol. 31, 844–848, 1983.
- Zaki, K. A. and A. Atia, "Modes in dielectric-loaded waveguides and resonators," *IEEE Trans. Microwave Theory Tech.*, Vol. 31, 1039–1045, 1983.
- Sheen, J., "Microwave measurements of dielectric properties using a closed cylindrical cavity dielectric resonator," *IEEE DEI*, Vol. 14, No. 15, 1139–1144, 2007.
- Sheen, J., C. A. Chen, Y. H. Chen, et al., "Microwave measurements of dielectric properties A further study to a new theoretical model for a closed cylindrical cavity dielectric resonator," *IEEE DEI*, Vol. 14, No. 12, 3874–3877, 2007.
- 22. Cao, L. and D. Cao, "Study of parameter Dielectric measurement of microwave dielectric materials by the closed cavity method," *Electronic Components and Materials*, Vol. 30, No. 12, 9–12, 2011, in Chinese.
- Zhou, D., M. Hu, S. Jiang, et al., "Microwave measurement of dielectric properties of ceramics by the closed cavity resonator method," J. Huazhong University of Sci. & Tech. (Nature Science Edition), Vol. 32, No. 8, 50–53, 2004, in Chinese.
- 24. Ni, E., *Measurement of Microwave Dielectric Resonator*, People's Posts and Telecommunications Press, Beijing, China, 2006.
- 25. Gu, J., *Dielectric Resonator Microwave Circuit*, People's Posts and Telecommunications Press, Beijing, China, 1986.